

**Problem Set 4**

1. Griffiths 12.6.
2. Griffiths 12.18.
3. Griffiths 12.19.
4. Griffiths 12.20.
5. Griffiths 12.32.

6. Inertial reference frames  $\mathcal{S}'$  and  $\mathcal{S}$  coincide at  $t' = t = 0$ . You may ignore the  $z$  dimension, so that a point in spacetime is determined by only three quantities  $r \equiv (ct, x, y)$ . The Lorentz transformation between  $\mathcal{S}$  and  $\mathcal{S}'$  is given by

$$\begin{pmatrix} ct' \\ x' \\ y' \end{pmatrix} = \mathcal{L} \begin{pmatrix} ct \\ x \\ y \end{pmatrix},$$

where  $\mathcal{L}$  is a  $3 \times 3$  matrix.

(a) Assume for this part that  $\mathcal{S}'$  moves with velocity

$$\mathbf{V} = \beta c \hat{\mathbf{x}}$$

with respect to  $\mathcal{S}$ . Using your knowledge of Lorentz transformations (no derivation necessary), write  $\mathcal{L}$  for this case.

(b) Assume for this part that  $\mathcal{S}'$  moves with velocity

$$\mathbf{V} = \beta c \frac{\hat{\mathbf{x}} + \hat{\mathbf{y}}}{\sqrt{2}}$$

with respect to  $\mathcal{S}$ . Find  $\mathcal{L}$  for this case. (*Hint.* Rotate to a system in which  $\mathbf{V}$  is along the  $\hat{\mathbf{x}}$  axis, transform using your answer for part (a.), and then rotate back. Check that your result is symmetric under interchange of  $x$  and  $y$ , as is  $\mathbf{V}$ , and that it reduces to the unit matrix as  $\beta \rightarrow 0$ .)

7. The now retired Bevatron at Berkeley Lab is famous for having produced the first observed antiprotons (you may have glimpsed white-maned Nobelist Owen Chamberlain, one of the first observers, being helped to his seat at

Physics Department colloquia). An economical reaction for producing antiprotons is

$$p + p \rightarrow p + p + p + \bar{p},$$

where the first proton is part of a beam, the second is at rest in a target, and  $\bar{p}$  is an antiproton. Because of the *CPT* theorem, both  $p$  and  $\bar{p}$  must have the same mass ( $= 0.94 \times 10^9$  eV).

At threshold, all four final state particles have essentially zero velocity *with respect to each other*. What is the beam energy in that case? (The actual Bevatron beam energy was  $6 \times 10^9$  eV).

8. (Taylor and Wheeler problem 51)

*The clock paradox, version 3.*

Can one go to a point 7000 light years away – and return – without aging more than 40 years? “Yes” is the conclusion reached by an engineer on the staff of a large aviation firm in a recent report. In his analysis the traveler experiences a constant “1 *g*” acceleration (or deceleration, depending on the stage reached in her journey). Assuming this limitation, is the engineer right in his conclusion? (For simplicity, limit attention to the first phase of the motion, during which the astronaut accelerates for 10 years – then double the distance covered in that time to find how far it is to the most remote point reached in the course of the journey.)

(a)

The acceleration is *not*  $g = 9.8$  meters per second per second relative to the laboratory frame. If it were, how many times faster than light would the spaceship be moving at the end of ten years (1 year =  $31.6 \times 10^6$  seconds)? *If the acceleration is not specified with respect to the laboratory, then with respect to what is it specified?* Discussion: Look at the bathroom scales on which one is standing! The rocket jet is always turned up to the point where these scales read one’s *correct* weight. Under these conditions one is being accelerated at 9.8 meters per second per second with respect to a spaceship

that (1) instantaneously happens to be riding alongside with identical velocity, but (2) is *not* being accelerated, and, therefore (3) *provides the (momentary) inertial frame of reference relative to which the acceleration is  $g$ .*

(b)

*How much velocity does the spaceship have after a given time?* This is the moment to object to the question and to rephrase it. *Velocity  $\beta c$*  is not the simple quantity to analyze. The simple quantity is the *boost parameter  $\eta$* . This parameter is simple because it is *additive* in this sense: Let the boost parameter of the spaceship with respect to the imaginary instantaneously comoving inertial frame change from 0 to  $d\eta$  in an astronaut time  $d\tau$ . Then the boost parameter of the spaceship with respect to the *laboratory* frame changes in the same astronaut time from its initial value  $\eta$  to the subsequent value  $\eta + d\eta$ . Now relate  $d\eta$  to the acceleration  $g$  in the instantaneously comoving inertial frame. In this frame  $g d\tau = c d\beta = c d(\tanh \eta) = (c/\cosh^2(\eta \approx 0)) d\eta \approx c d\eta$  so that

$$c d\eta = g d\tau$$

Each lapse of time  $d\tau$  on the astronaut's watch is accompanied by an additional increase  $d\eta = \frac{g}{c} d\tau$  in the boost parameter of the spaceship. In the laboratory frame the total boost parameter of the spaceship is simply the sum of these additional increases in the boost parameter. Assume that the spaceship starts from rest. Then its boost parameter will increase linearly with *astronaut* time according to the equation

$$c\eta = g\tau$$

This expression gives the boost parameter  $\eta$  of the spaceship in the *laboratory* frame at any time  $\tau$  in the *astronaut's* frame.

(c)

*What laboratory distance  $x$  does the spaceship cover in a given astronaut time  $\tau$ ?* At any instant the velocity of the spaceship in the laboratory frame is related to its boost parameter by the equation  $dx/dt = c \tanh \eta$  so that the distance  $dx$  covered in *laboratory* time  $dt$  is

$$dx = c \tanh \eta dt$$

Remember that the time between ticks of the astronaut's watch  $d\tau$  appear to have the larger value  $dt$  in the laboratory frame (time dilation) given by the expression

$$dt = \cosh \eta d\tau$$

Hence the laboratory distance  $dx$  covered in *astronaut* time  $d\tau$  is

$$dx = c \tanh \eta \cosh \eta d\tau = c \sinh \eta d\tau$$

Use the expression  $c\eta = g\tau$  from part b to obtain

$$dx = c \sinh \left( \frac{g\tau}{c} \right) d\tau$$

Sum (integrate) all these small displacements  $dx$  from zero astronaut time to a final astronaut time to find

$$x = \frac{c^2}{g} \left[ \cosh \left( \frac{g\tau}{c} \right) - 1 \right]$$

This expression gives the laboratory *distance  $x$*  covered by the spaceship at any time  $\tau$  in the astronaut's frame.

(d)

Plugging in the appropriate numerical values, determine whether the engineer is correct in his conclusion reported at the beginning of this exercise.